



Institute of  
Continuum Mechanics

Leibniz  
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Hannover

Seventh Meeting  
UNILATERAL PROBLEMS IN  
STRUCTURAL ANALYSIS

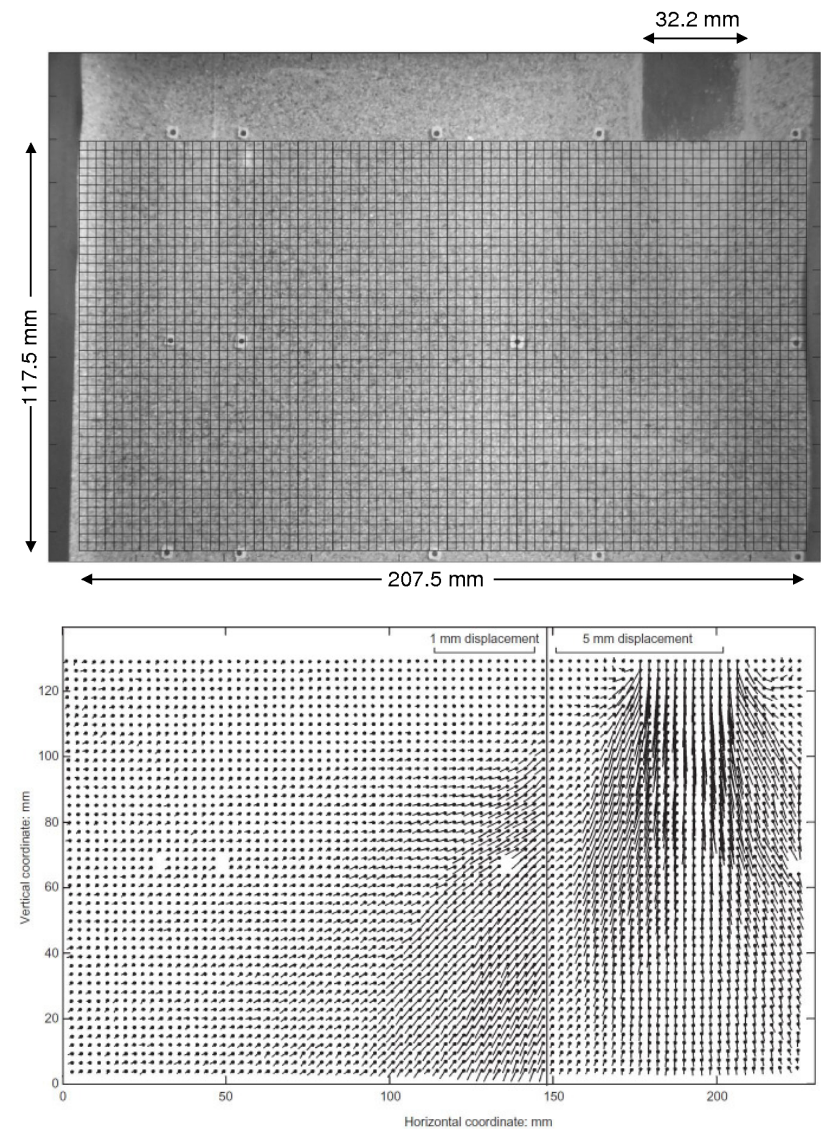
June 17-19 2010  
Palmanova, Italy

# Multiscale Modeling of Non-Cohesive Granular Materials

Christian Wellmann, Peter Wriggers

## motivation

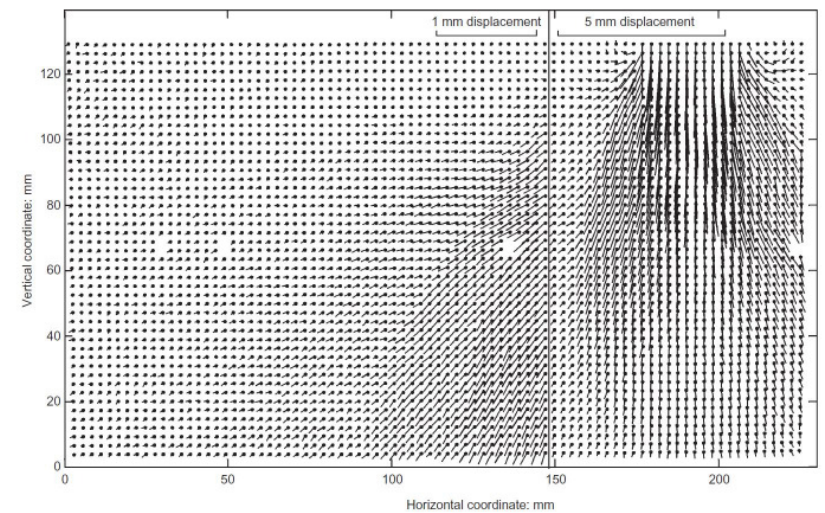
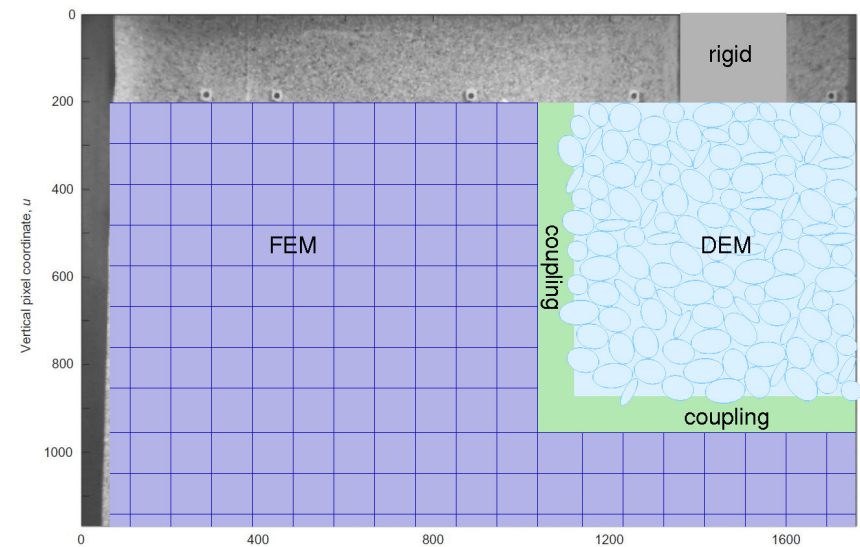
- complex material behavior
- phenomenological approach
  - assume continuum
  - elasto-plastic constitutive equation
- drawbacks
  - numerous parameters
  - stability, accuracy and technical issues for large deformations
- multi-scale model
  - particle model for “hot-spot” regions
  - continuum model elsewhere
- requires coupling at interface
- fit continuum model to homogenized particle behavior



pictures from [White & Bolton, 2004]

## motivation

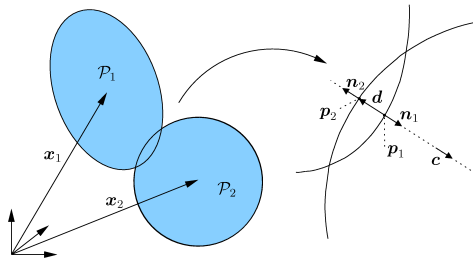
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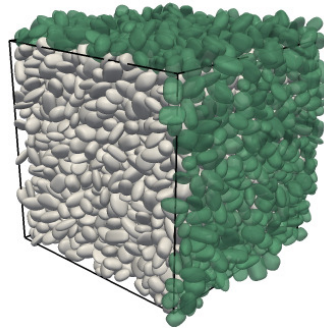
pictures from [White & Bolton, 2004]

outline

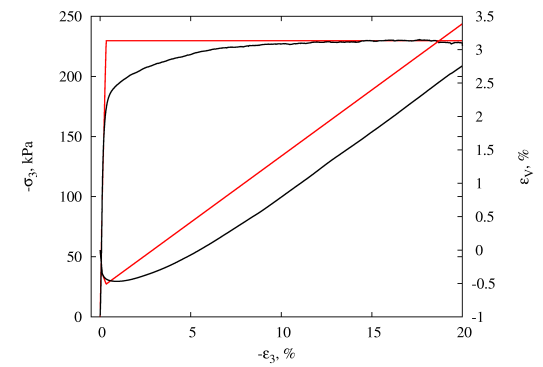
discrete element model



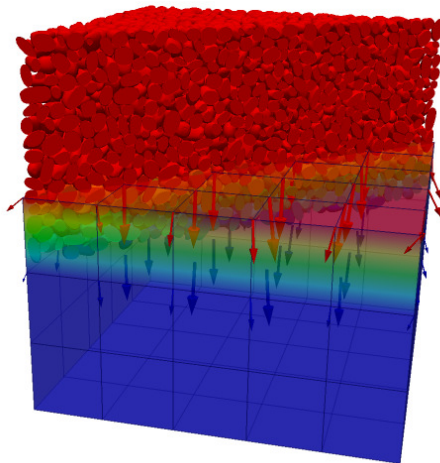
homogenization



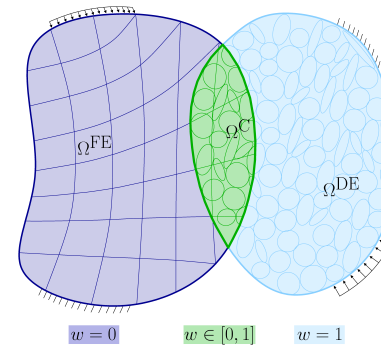
fit elasto-plastic parameters



numerical examples



DE-FE coupling scheme



## discrete element method

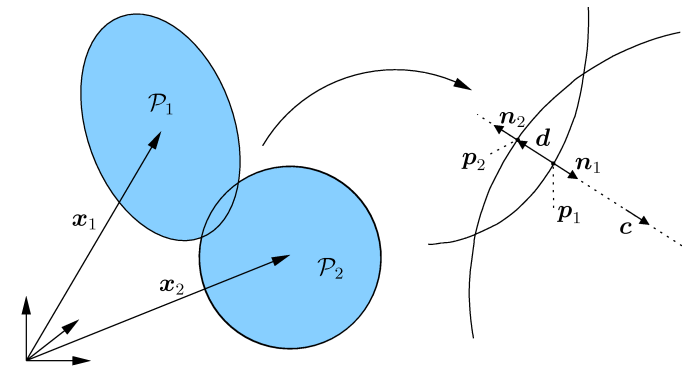
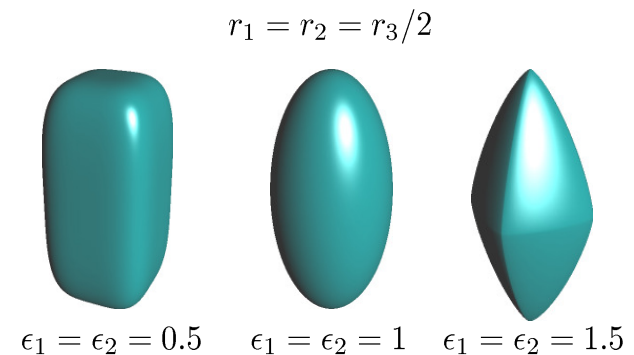
- dry non-cohesive frictional granular material
- individual grains are modeled as rigid bodies
- superquadric geometry [Barr, 1981]

$$F(\mathbf{x}) = \left( \left| \frac{x_1}{r_1} \right|^{\frac{2}{\epsilon_1}} + \left| \frac{x_2}{r_2} \right|^{\frac{2}{\epsilon_1}} \right)^{\frac{\epsilon_1}{\epsilon_2}} + \left| \frac{x_3}{r_3} \right|^{\frac{2}{\epsilon_2}} = 1$$

- arbitrary aspect ratios & angularity
- contact:
  - admit small overlap
  - Hertz-Mindlin contact theory
  - Coulomb friction
- explicit time integration

$$m \ddot{\mathbf{x}} = \mathbf{f}^{\text{res}}$$

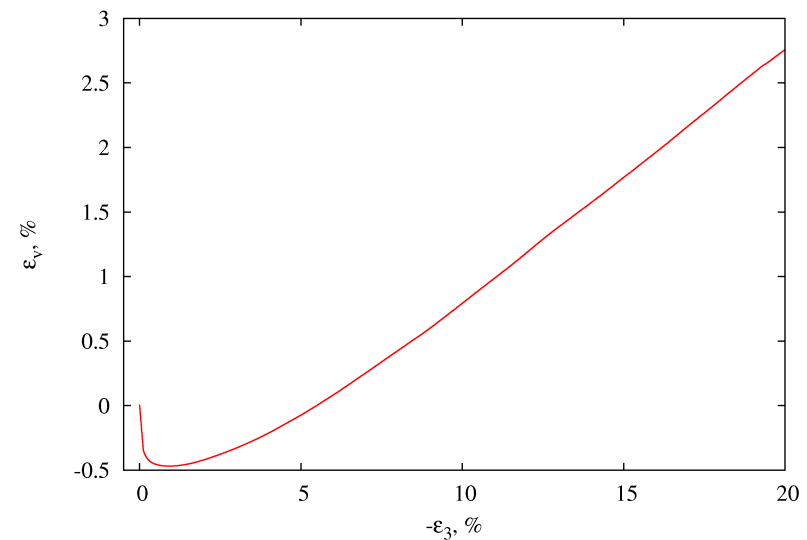
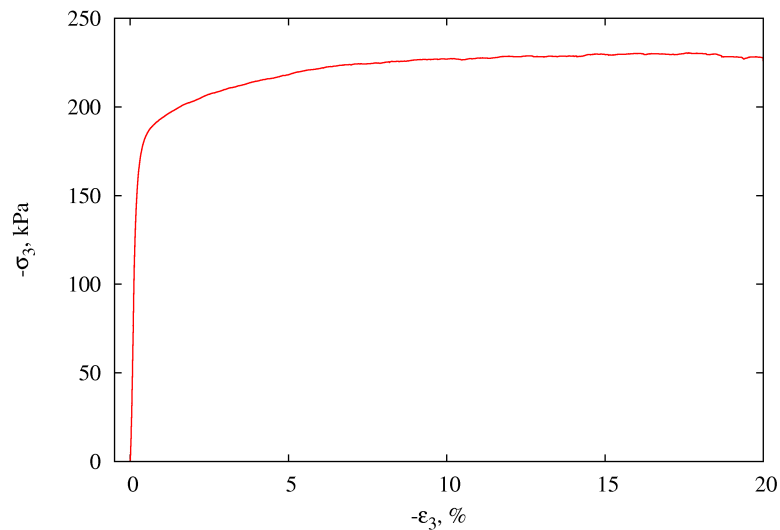
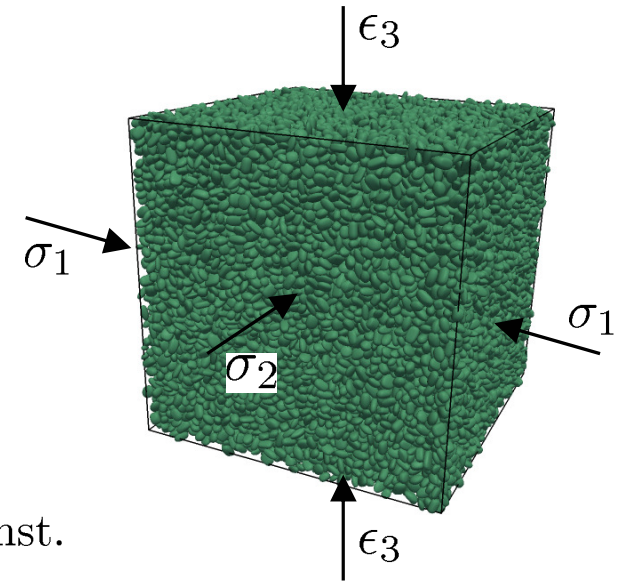
$$\mathbf{I} \cdot \dot{\boldsymbol{\omega}} + \boldsymbol{\omega} \times \mathbf{I} \cdot \boldsymbol{\omega} = \mathbf{m}^{\text{res}}$$



## homogenization

- random periodic unit cell
- DEM parameters adapted to Leighton Buzzard Sand fraction B
  - size: 0.6 – 1.18 mm
  - $E=50$  GPa,  $\nu=0.2$
  - friction [Rowe, 1962], [Ishibashi, 1994]:  $\mu=0.24$
- triaxial test with adaptive dimension control:
 
$$\sigma_1 = \sigma_2 = \text{const.}$$

$$\dot{\epsilon}_3 = \text{const.}$$





## fitting of constitutive equation

- non-associative Mohr-Coulomb

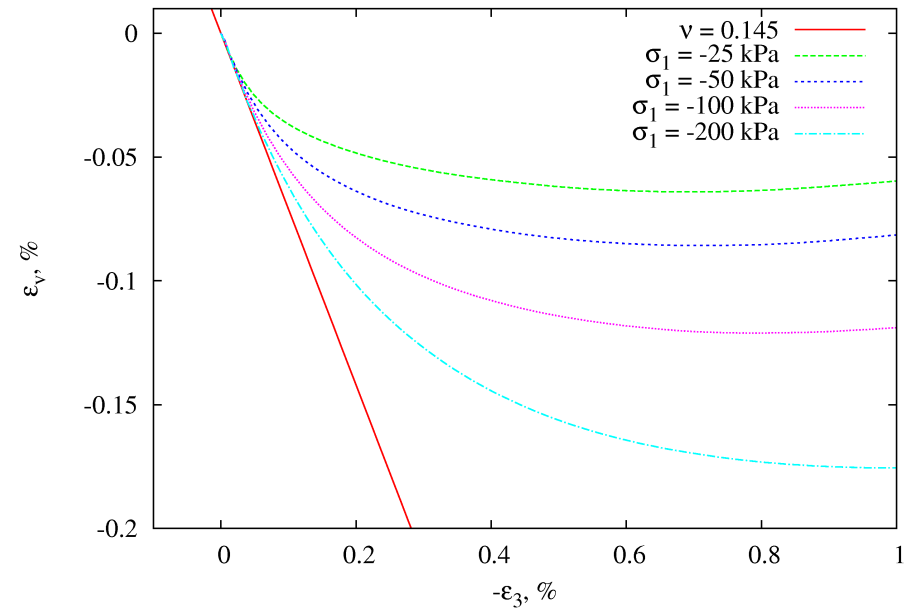
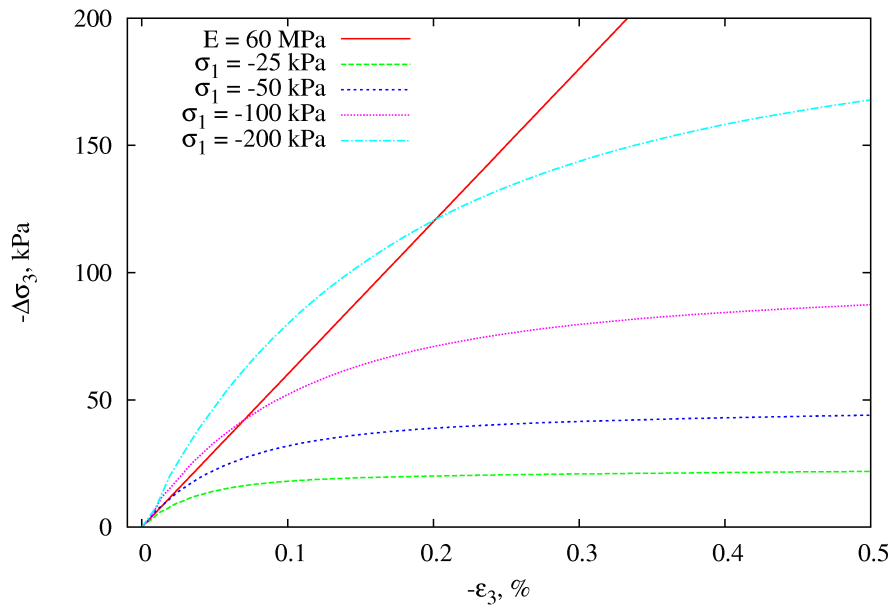
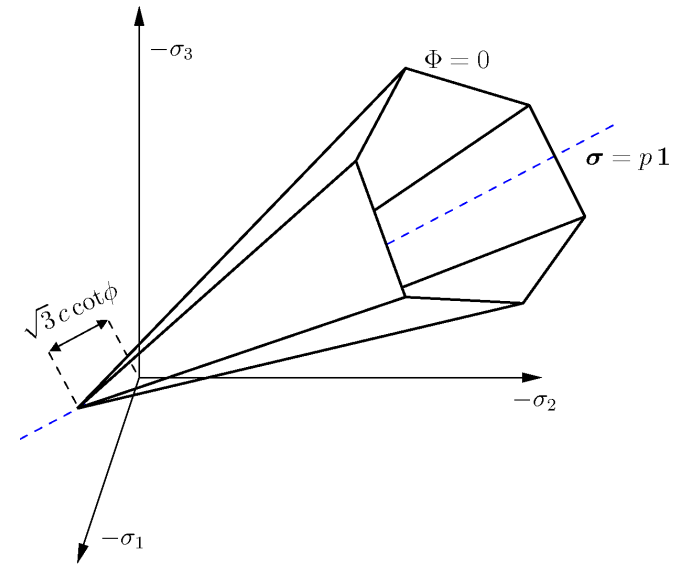
$$\epsilon = \epsilon^e + \epsilon^p$$

$$\sigma = 2G \epsilon_{\text{dev}}^e + K \epsilon_V^e \mathbf{1}$$

$$\Phi = \sigma_1 - \sigma_3 + (\sigma_1 + \sigma_3) \sin \varphi + 2c \cos \varphi$$

$$\Psi = \sigma_1 - \sigma_3 + (\sigma_1 + \sigma_3) \sin \psi + 2c \cos \psi$$

$$\Rightarrow G, K, \varphi, \psi, (c)$$





## fitting of constitutive equation

- non-associative Mohr-Coulomb

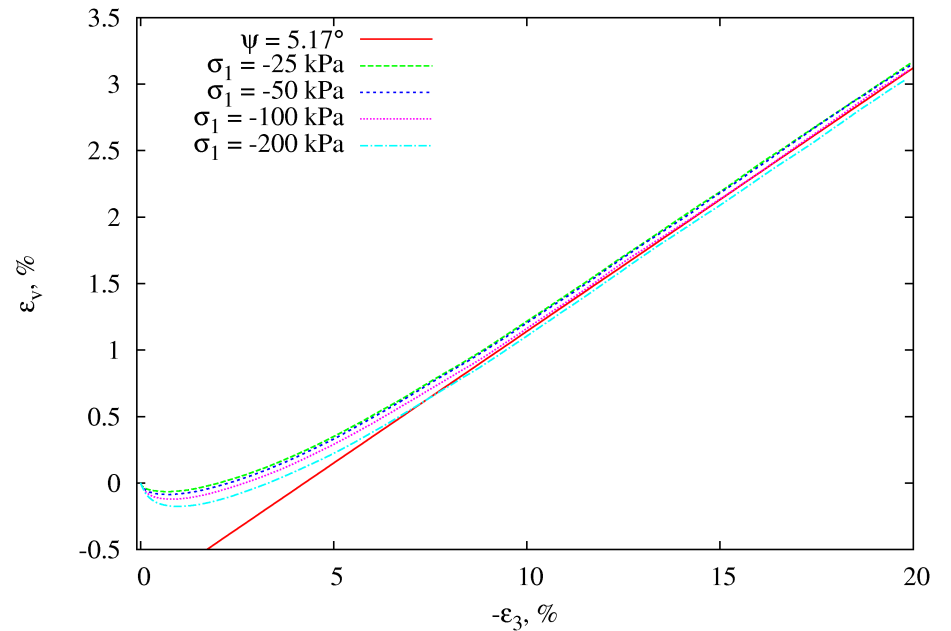
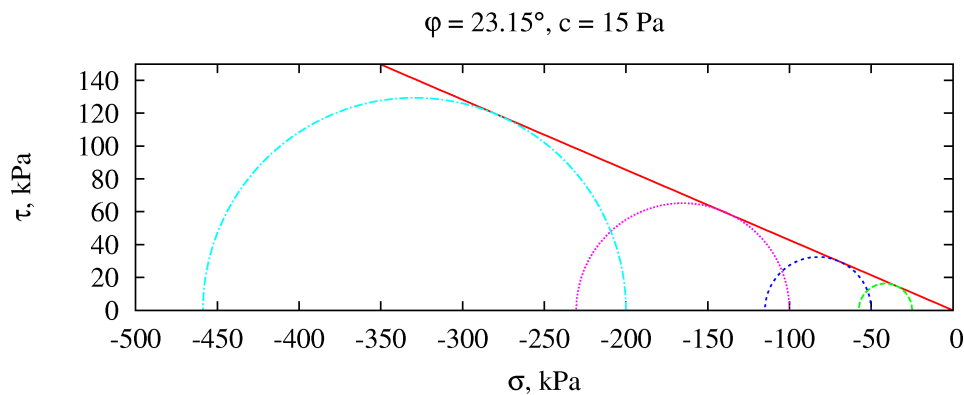
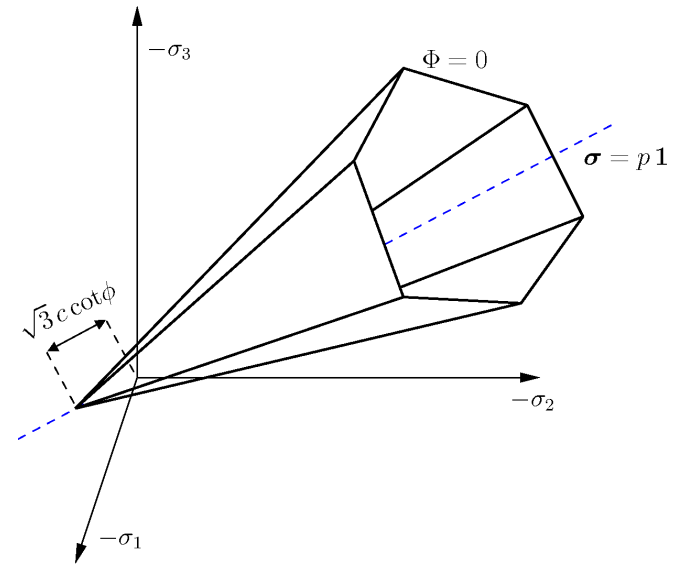
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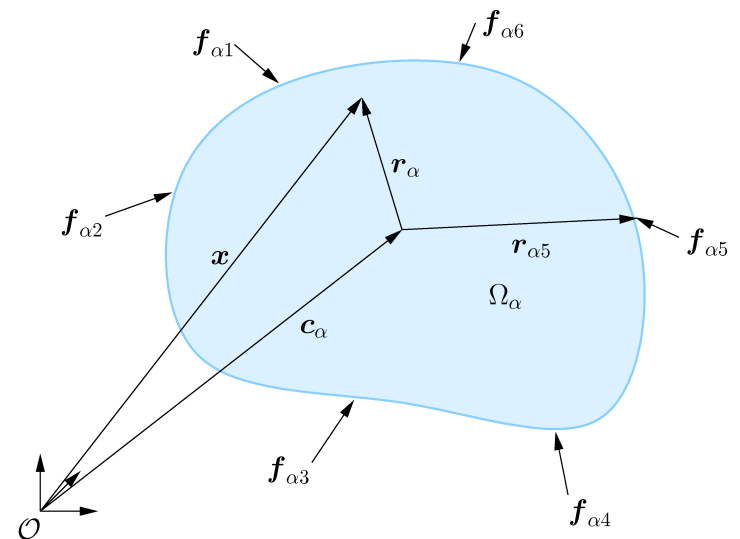
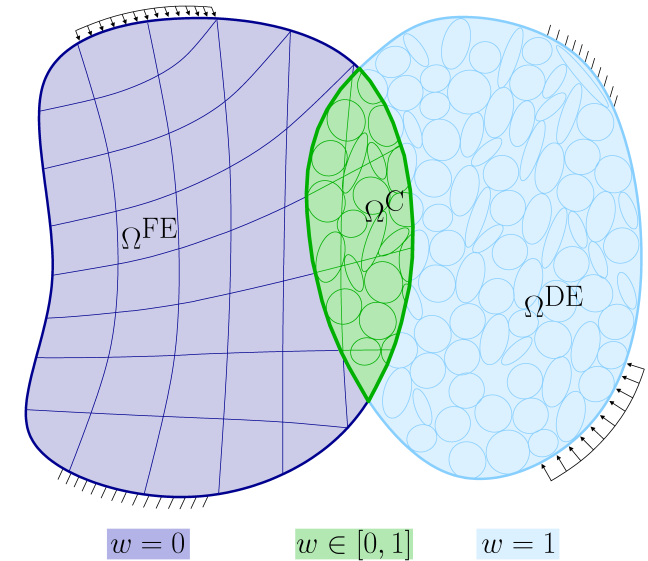


## coupling discrete & finite element method

- use Arlequin method [Ben Dhia, 1998]
  - both models overlap in coupling domain  $\Omega^C$
  - introduce weight factor to interpolate virtual work

$$\begin{aligned} \delta W &= \delta W^{\text{FE}} + \delta W^{\text{DE}} \quad \text{with} \\ \delta W^{\text{FE}} &= \int_{\Omega^{\text{FE}}} (1-w) \boldsymbol{\sigma} : \delta \boldsymbol{\epsilon} \, d\Omega + \int_{\Omega^{\text{FE}}} (1-w) \rho (\ddot{\mathbf{x}} - \mathbf{b}) \cdot \delta \mathbf{u} \, d\Omega - \\ &\quad \int_{\Gamma^{\text{FE}}} (1-w) \mathbf{t} \cdot \delta \mathbf{u} \, d\Gamma \\ \delta W^{\text{DE}} &= \sum_{\alpha=1}^{n_p} \delta W_{\alpha} \\ &= \sum_{\alpha=1}^{n_p} \left[ \int_{\Omega_{\alpha}} w \rho (\ddot{\mathbf{x}} - \mathbf{b}) \cdot \delta \mathbf{u}_{\alpha} \, d\Omega - \sum_{\beta=1}^{n_{\alpha}} w_{\alpha\beta} \mathbf{f}_{\alpha\beta} \cdot \delta \mathbf{u}_{\alpha} \right] \end{aligned}$$

- FE part: standard integration with additional weight factor





## weighted virtual work of a rigid body

$$\delta W = \int_{\Omega} w \rho (\ddot{\mathbf{x}} - \mathbf{b}) \cdot \delta \mathbf{u} \, d\Omega - \sum_{\beta=1}^n w_{\beta} \mathbf{f}_{\beta} \cdot \delta \mathbf{u}$$

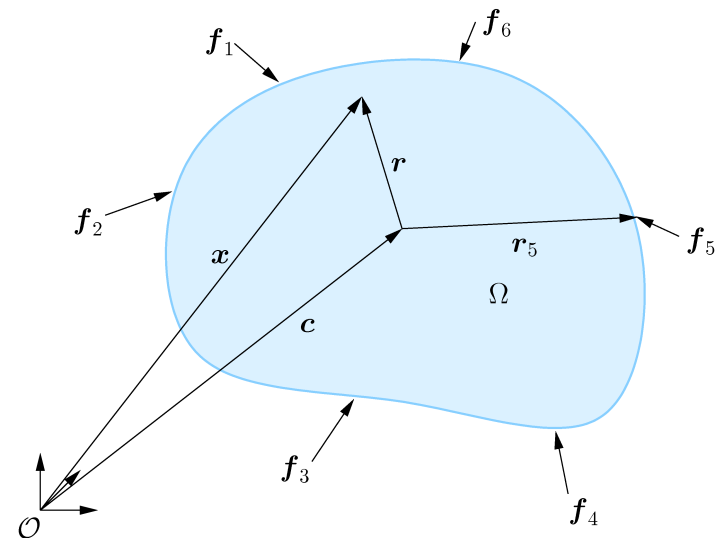
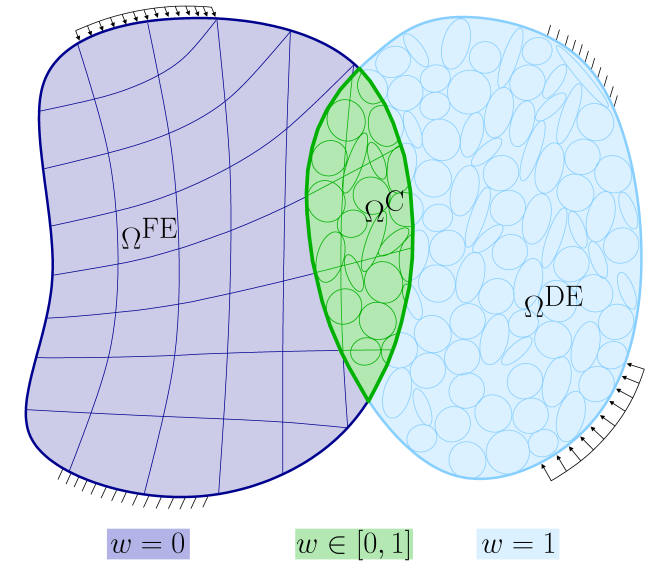
$$\mathbf{x} = \mathbf{c} + \mathbf{r} \Rightarrow \ddot{\mathbf{x}} = \ddot{\mathbf{c}} + \dot{\boldsymbol{\omega}} \times \mathbf{r} + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r})$$

$$\delta \mathbf{u}(\mathbf{x}) = \delta \mathbf{u}^0 + \delta \boldsymbol{\omega} \times \mathbf{r}$$

- assuming  $w(\mathbf{x})$  smooth, continuous, monotonic

$$w(\mathbf{x}) \approx w(\mathbf{c}) + \text{grad } w|_{\mathbf{c}} \cdot \mathbf{r} = w_c + w'_c \mathbf{n} \cdot \mathbf{r}, \quad \|\mathbf{n}\| = 1$$

$$\delta W = \left[ w_c m (\ddot{\mathbf{c}} - \mathbf{b}) - \sum_{\beta=1}^n w_{\beta} \mathbf{f}_{\beta} \right] \cdot \delta \mathbf{u}^0 + \left[ w_c (\mathbf{I} \cdot \dot{\boldsymbol{\omega}} + \boldsymbol{\omega} \times \mathbf{I} \cdot \boldsymbol{\omega}) - \sum_{\beta=1}^n w_{\beta} \mathbf{r}_{\beta} \times \mathbf{f}_{\beta} \right] \cdot \delta \boldsymbol{\omega} + \underbrace{w'_c \mathbf{n} \cdot \int_{\Omega} \mathbf{r} \rho (\ddot{\mathbf{c}} + \dot{\boldsymbol{\omega}} \times \mathbf{r} + \boldsymbol{\omega} \times (\boldsymbol{\omega} \times \mathbf{r}) - \mathbf{b}) \cdot (\delta \mathbf{u}_0 + \delta \boldsymbol{\omega} \times \mathbf{r}) \, d\Omega}_{\text{negligible for } \Omega \ll \Omega^C}$$





## kinematic constraints

- discrete  $\rightarrow$  continuous projection:

$$\Pi(\mathbf{u}_\alpha, \mathbf{u}_\beta, \dots) = \mathbf{u}^{\text{DE}}(\mathbf{x}) \quad \text{with} \quad \alpha, \beta, \dots \in \mathcal{P}^C$$

- $L^2$  penalty constraint term

$$C = \frac{\epsilon}{2} \int_{\Omega^C} \|\mathbf{u}^{\text{DE}}(\mathbf{x}) - \mathbf{u}^{\text{FE}}(\mathbf{x})\|^2 d\Omega$$

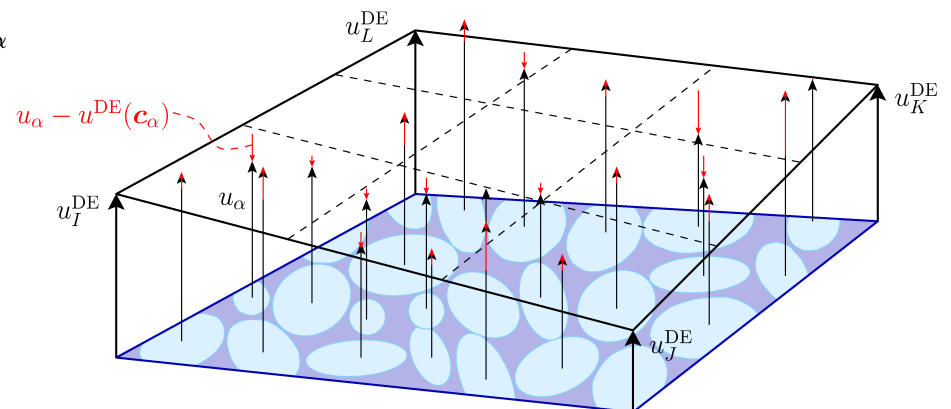
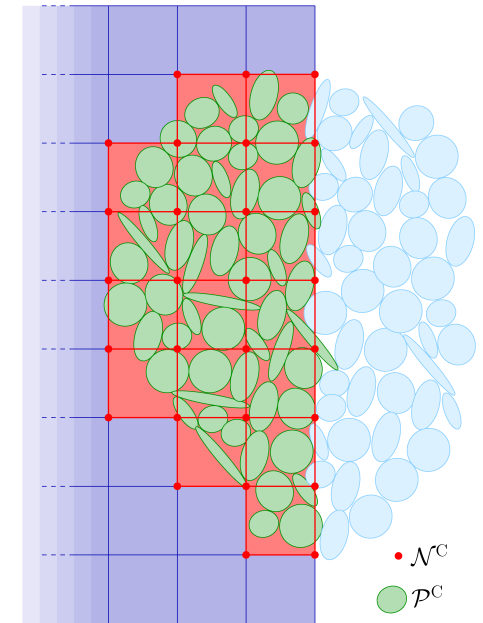
- projection via LSF using FE ansatz  
 $\Rightarrow$  split into coarse & fine scale [Wagner & Liu, 2003]

$$\mathbf{u}^{\text{DE}}(\mathbf{x}) = \sum_{I \in \mathcal{N}^C} N_I \mathbf{u}_I^{\text{DE}}, \quad \min_{\mathbf{u}_I^{\text{DE}}} \sum_{\alpha \in \mathcal{P}^C} V_\alpha \|\mathbf{u}_\alpha - \mathbf{u}^{\text{DE}}(\mathbf{c}_\alpha)\|^2$$

$$\Rightarrow \sum_{\alpha \in \mathcal{P}^C} \sum_{J \in \mathcal{N}^C} N_{I\alpha} V_\alpha N_{J\alpha} \mathbf{u}_J^{\text{DE}} = \sum_{\alpha \in \mathcal{P}^C} N_{I\alpha} V_\alpha \mathbf{u}_\alpha$$

$$\Rightarrow \underbrace{\underline{N} \underline{V} \underline{N}^T}_{\underline{A}} \underline{u}_c = \underline{N} \underline{V} \underline{u}_d$$

$$\underline{u}_c = \underbrace{\underline{A}^{-1} \underline{N} \underline{V}}_{\underline{\Pi}} \underline{u}_d \quad \text{with} \quad N_{I\alpha} = N_I(\mathbf{c}_\alpha)$$





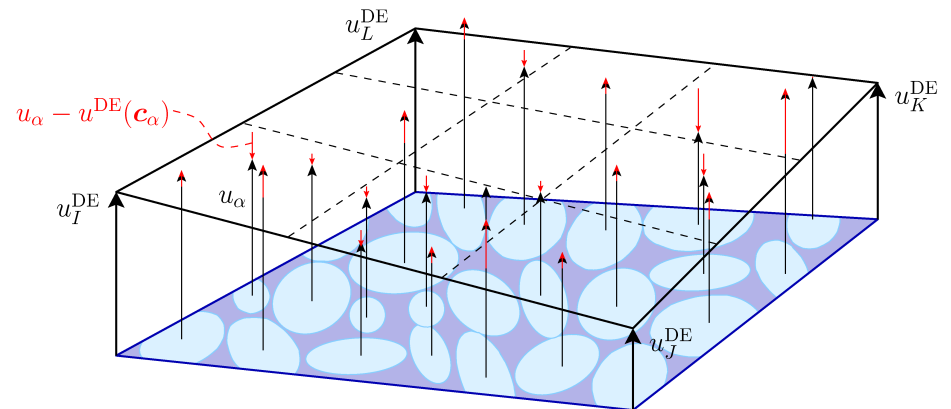
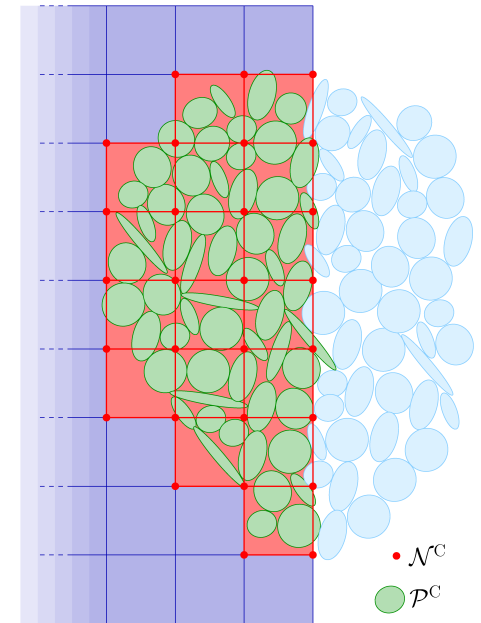
## coupling forces

- variational formulation with penalty term

$$\begin{aligned}
 \delta W + \delta C &= 0 \quad \text{with} \\
 \delta C &= \epsilon \int_{\Omega^c} \underbrace{(\mathbf{u}^{\text{DE}} - \mathbf{u}^{\text{FE}})}_{\mathbf{r}} \cdot (\delta \mathbf{u}^{\text{DE}} - \delta \mathbf{u}^{\text{FE}}) d\Omega \\
 &= \sum_{I \in \mathcal{N}^c} \delta \mathbf{r}_I \cdot \left( \epsilon \sum_{J \in \mathcal{N}^c} \int_{\Omega^c} N_I N_J d\Omega \mathbf{r}_J \right) \\
 &= \sum_{I \in \mathcal{N}^c} \delta \mathbf{r}_I \cdot \underbrace{\left( \epsilon \sum_{J \in \mathcal{N}^c} V_{IJ} \mathbf{r}_J \right)}_{\mathbf{f}_I^c}
 \end{aligned}$$

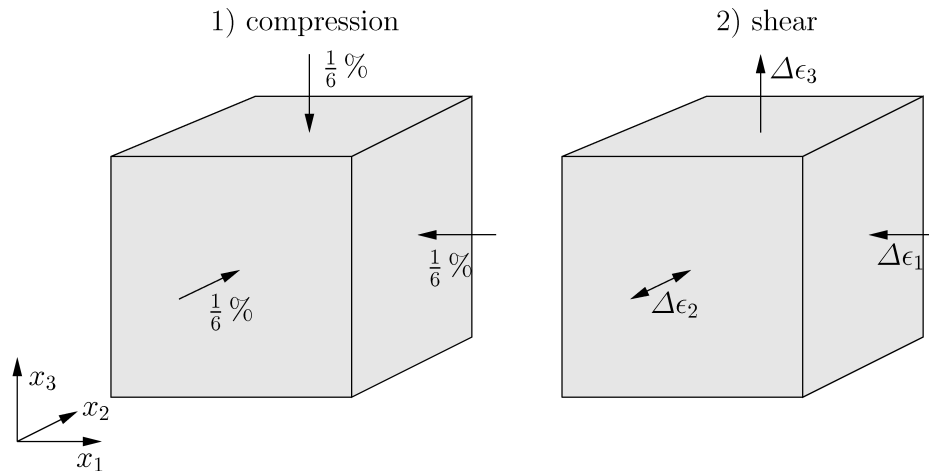
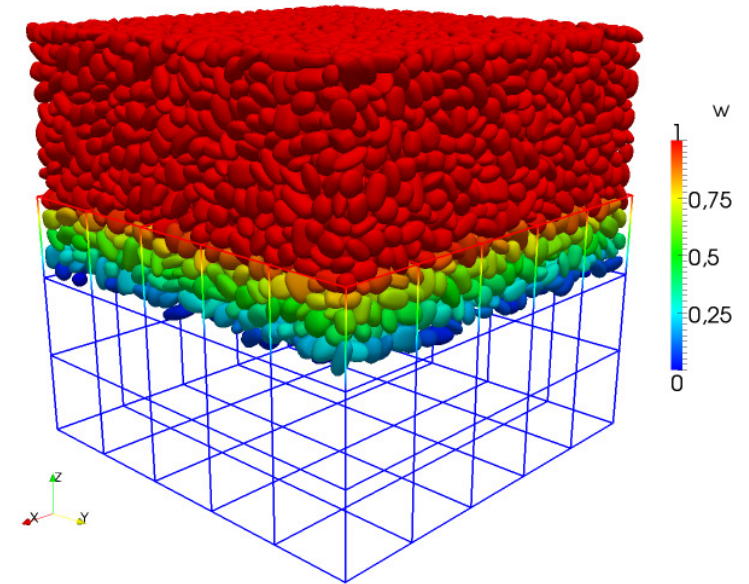
- the projection yields

$$\begin{aligned}
 \delta \mathbf{u}_I^{\text{DE}} &= \sum_{\alpha \in \mathcal{P}^c} \Pi_{I\alpha} \delta \mathbf{u}_\alpha^0 \\
 \Rightarrow \mathbf{f}_\alpha^c &= - \sum_{I \in \mathcal{N}^c} \Pi_{I\alpha} \mathbf{f}_I^c
 \end{aligned}$$



## numerical example: triaxial test

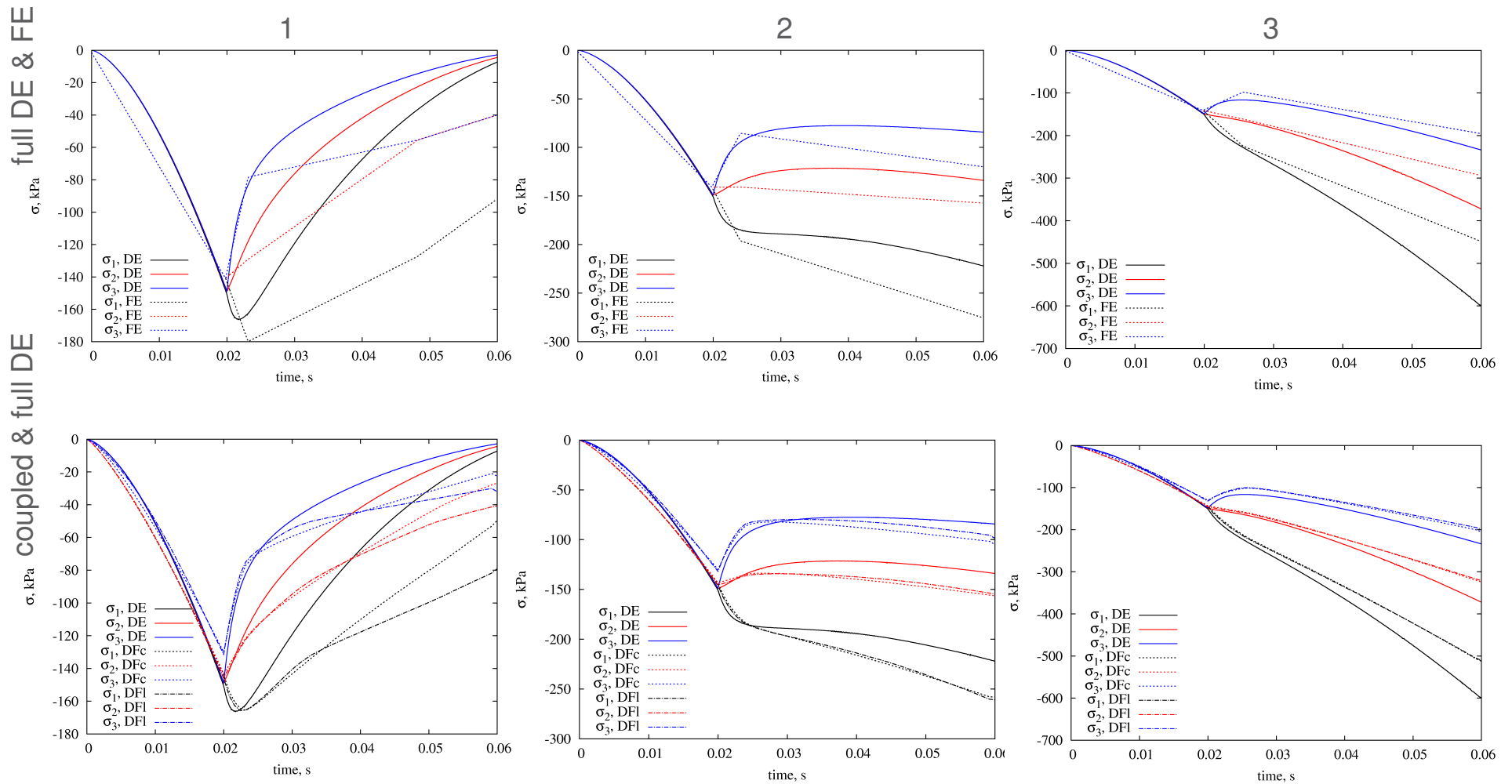
- system:
  - $\Omega^C = 1$  element layer
  - weight function defined on FE ansatz
  - trilinear hexahedral elements
  - full DE:  $\sim 14,000$  particles
  - coupled:  $\sim 8,000$  particles
  - $\epsilon = 10^8 \text{ Pa/mm}^2$
- loading:
  - 3 displacement controlled tests with two phases:



test	1	2	3
$\Delta\epsilon_V, \%$	1/2	0	-1/2
$\Delta\epsilon_1, \%$	-1	-1	-1
$\Delta\epsilon_2, \%$	1/6	0	-1/6
$\Delta\epsilon_3, \%$	4/3	1	2/3



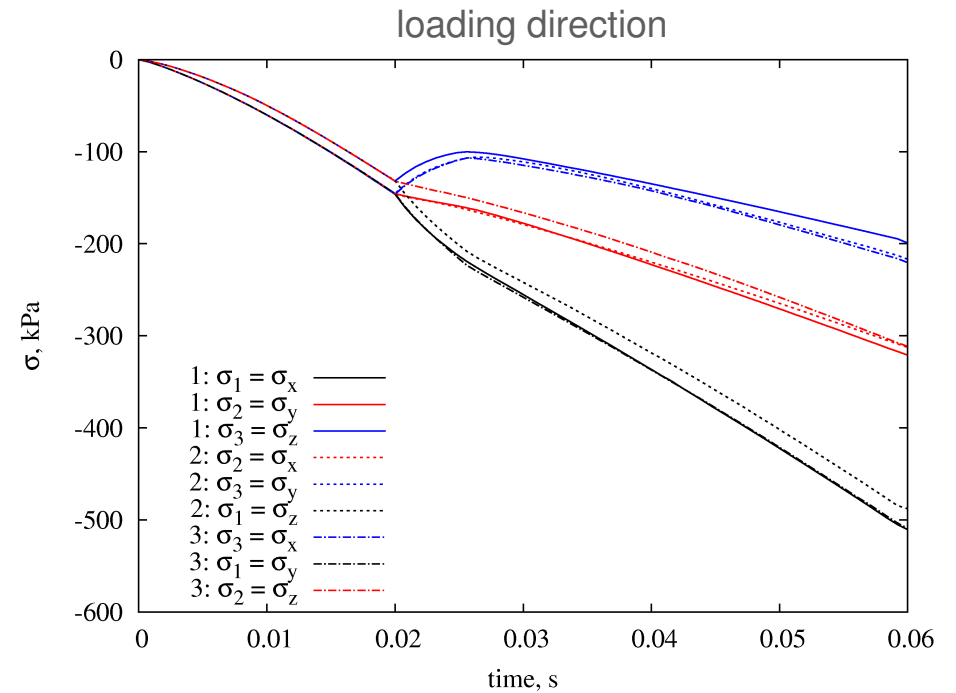
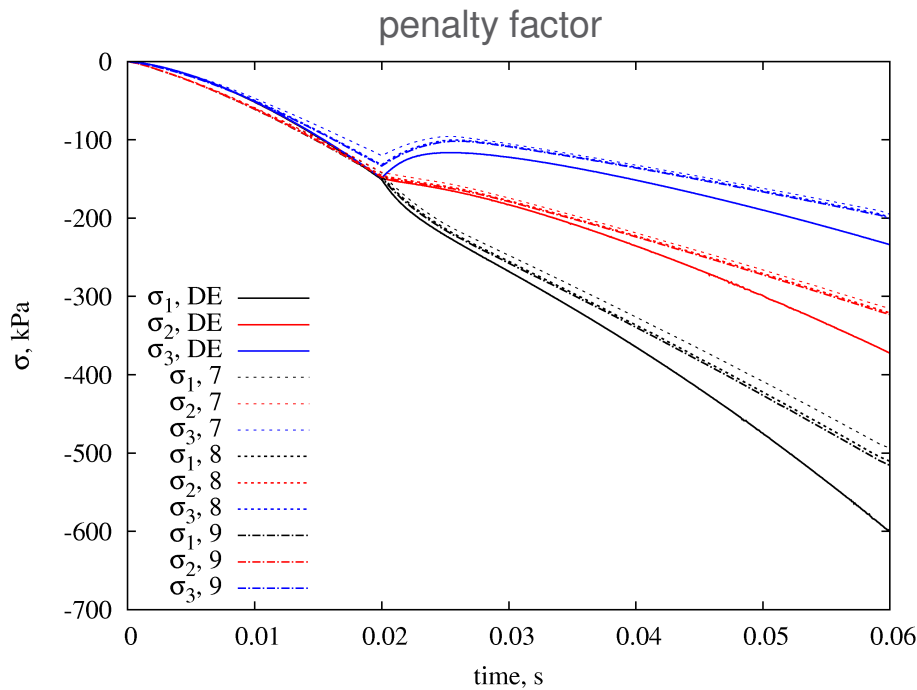
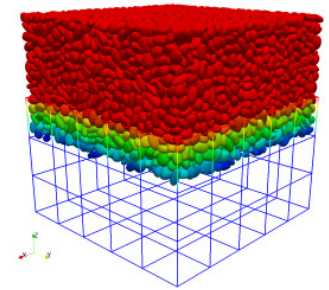
# results



## penalty factor & loading direction

- 3 penalty factors  
 $\epsilon = 10^7, 10^8, 10^9 \text{ Pa/mm}^2$
- $\Rightarrow 10^8$  is sufficient

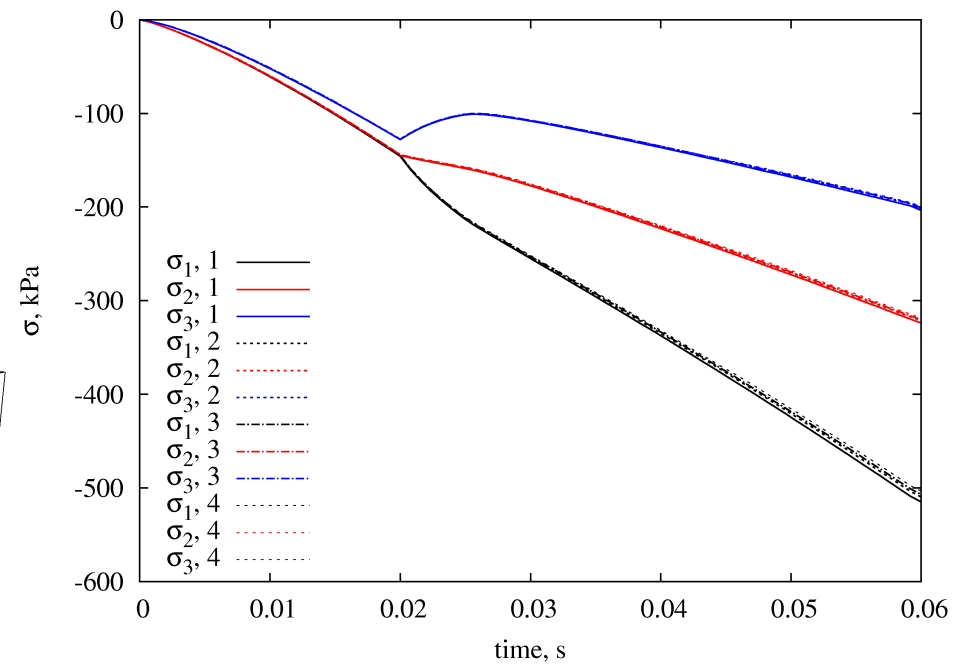
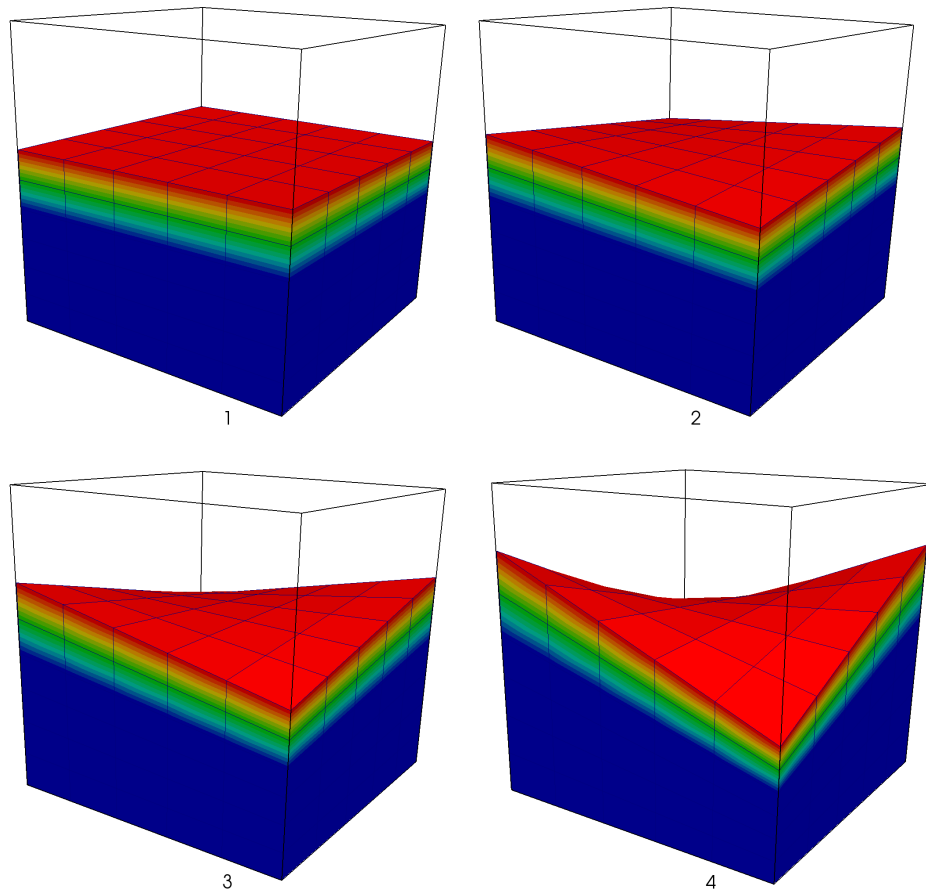
- switched loading directions
- small deviations in isotropic compression







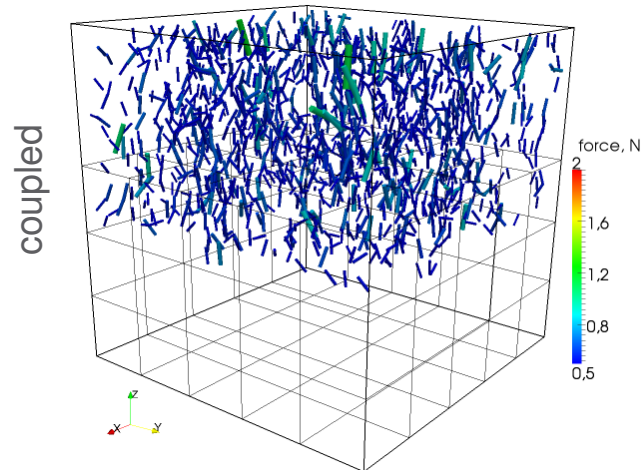
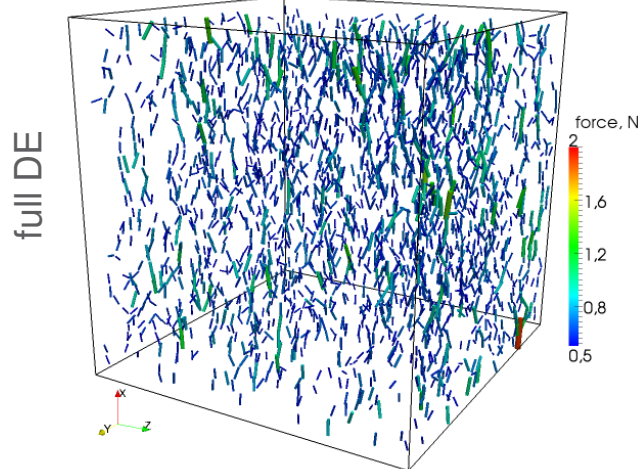
## coupling geometry



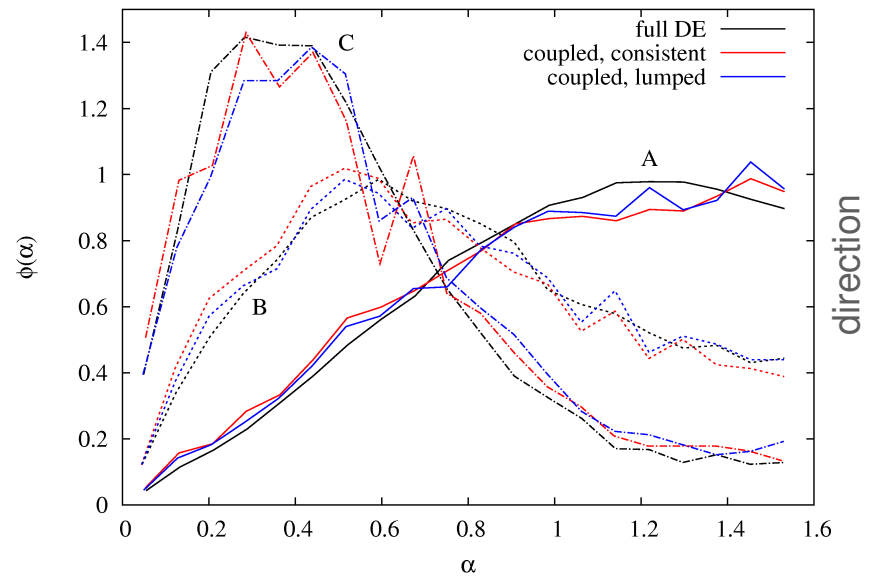
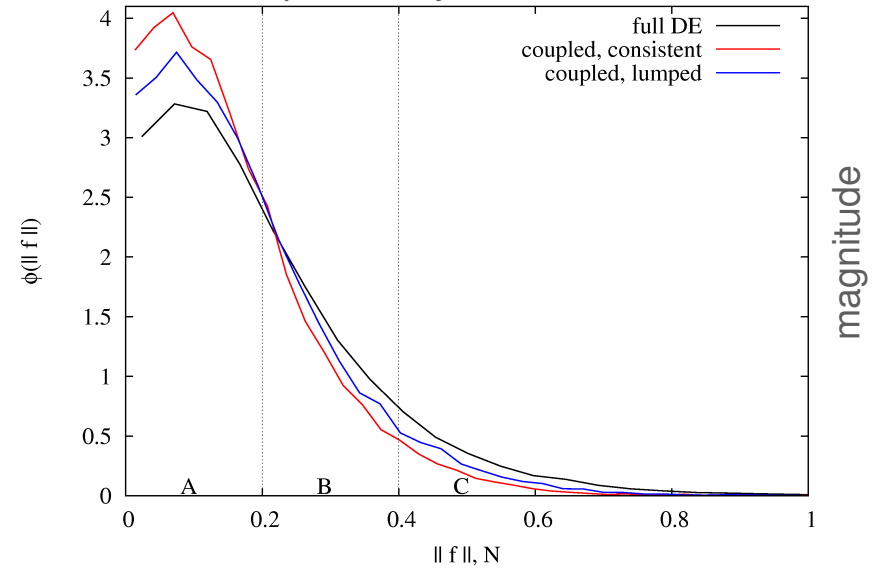


## microstructure

### contact force chains

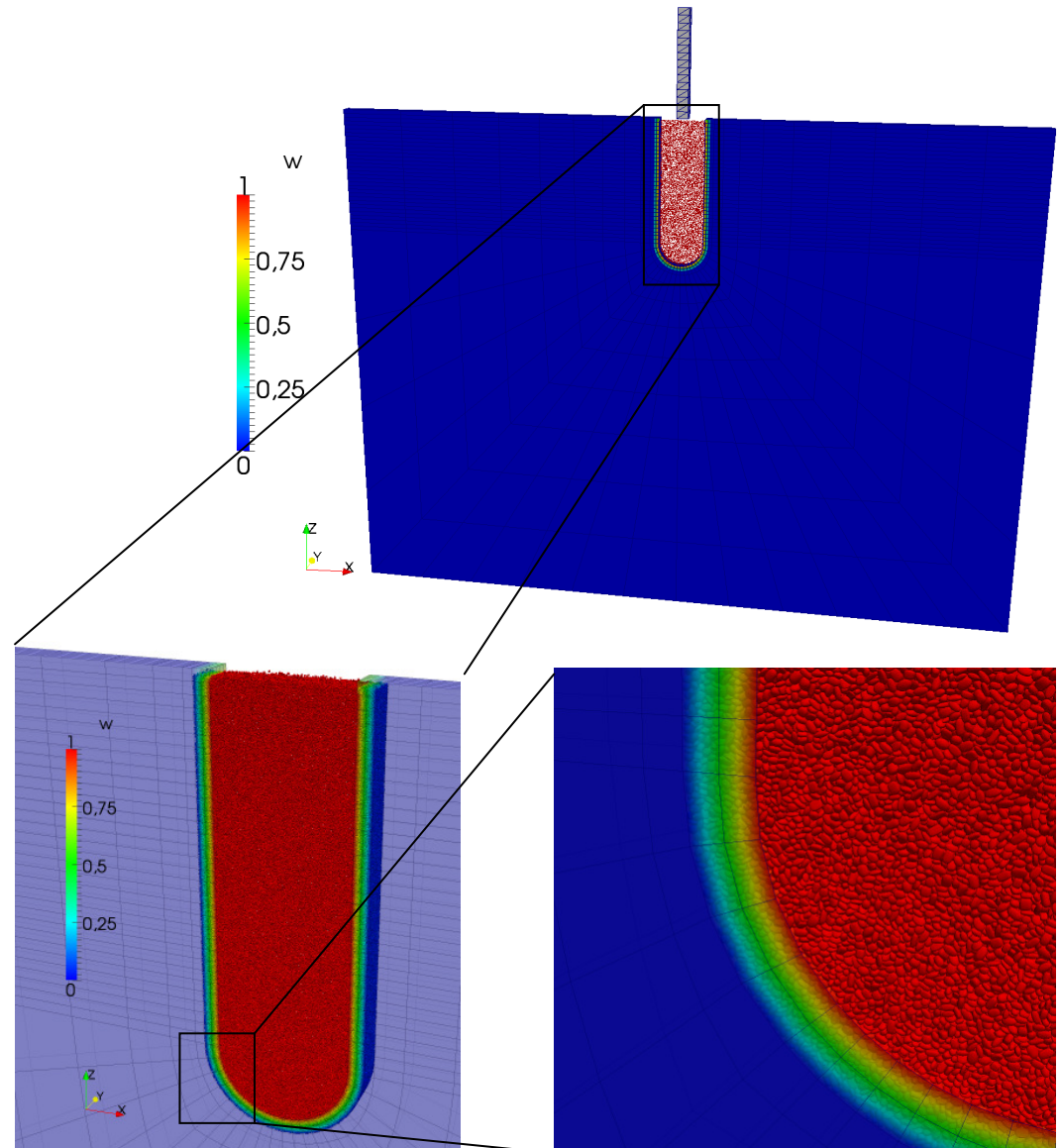


### probability distributions

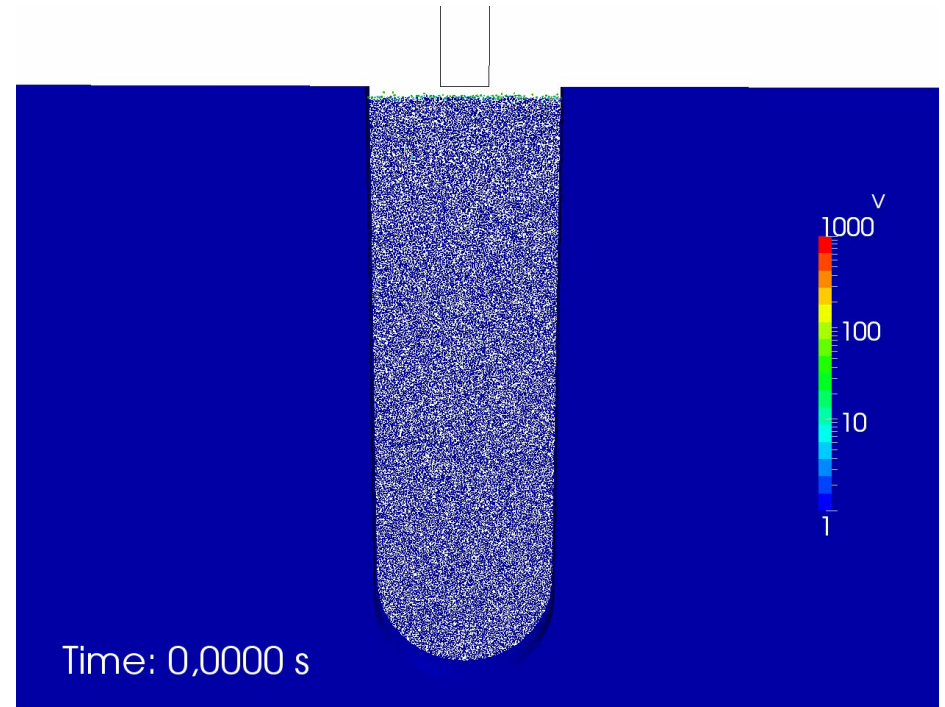
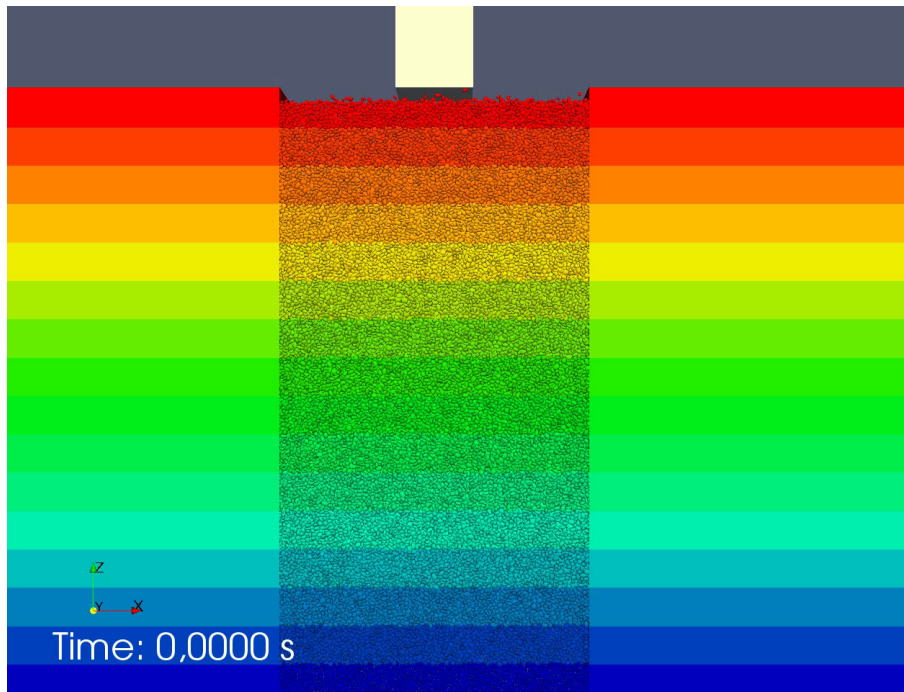


## pile installation

- box: 1000 x 750 x 16.1 mm
- pile cross section: 16.1 x 16.1 mm
- pile velocity: 1 m/s
- friction pile-particles: 0.1
- full DE model:  $\sim 25 \times 10^6$  particles
- large deformations only in pile vicinity
- coupled simulation
  - 540,000 particles
  - 5040 trilinear bricks
  - 7560 nodes
  - $3.2 \times 10^6$  DOFs
  - lumped projection



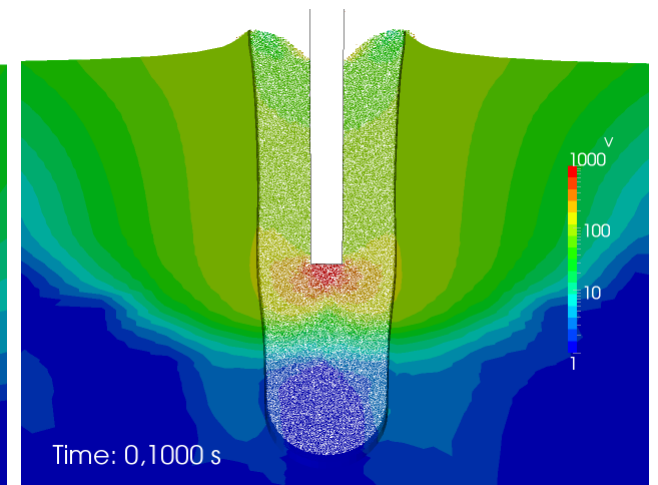
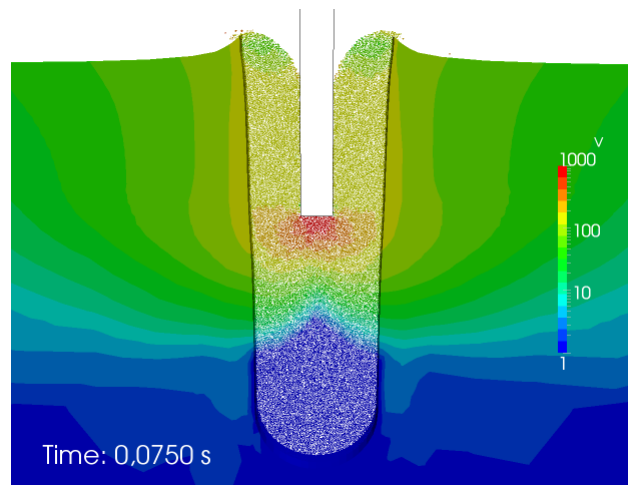
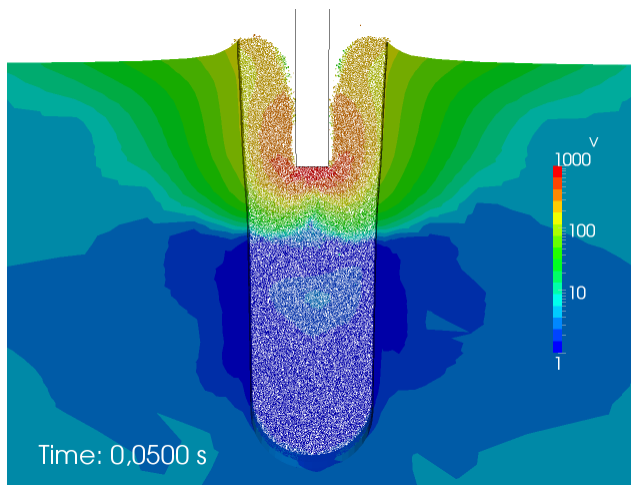
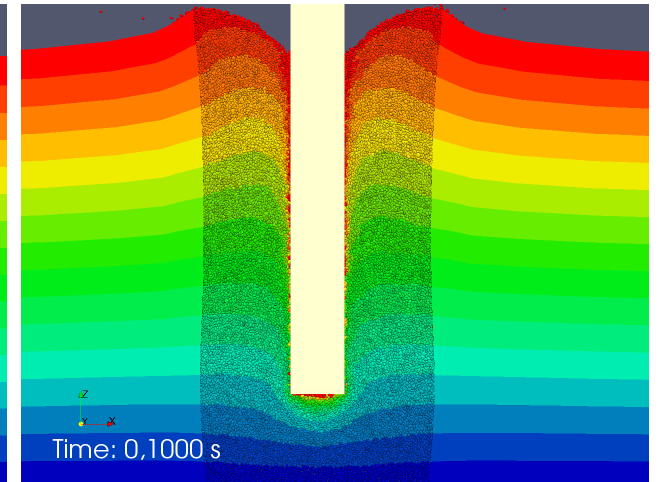
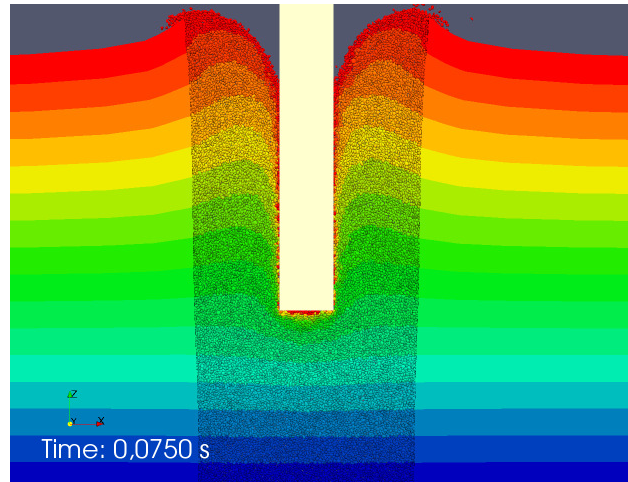
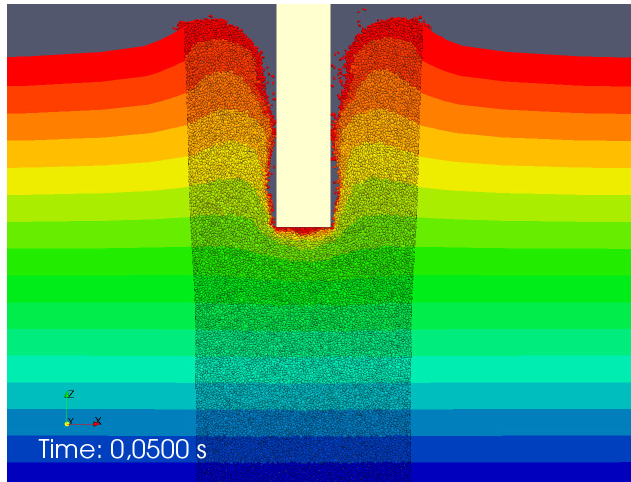
## results



- particle cone forms at flat pile tip
- material is “spread” by cone and moves outside and upwards
- smooth transition between DE & FE domain



## results



## summary & outlook

- DE-FE coupling scheme:
  - overlapping domain
  - weighted virtual work
  - coarse-fine split of DE displacements
- triaxial test:
  - similar results for consistent & lumped projection
  - works for distorted coupling geometry
  - preserves characteristic microstructure
- coupling scheme enables simulations not possible with mono-method model
- outlook:
  - enhance particle model to better represent shear strength of real granular materials
  - implement constitutive equation giving a better fit of the homogenized DE behavior
  - develop criteria and methods to adaptively control DE domain

